

QCD and Strings

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- A 35-years-old Puzzle
- AdS/CFT correspondence
- AdS/CFT and QCD Hydrodynamics
- Thermalization of a perfect fluid

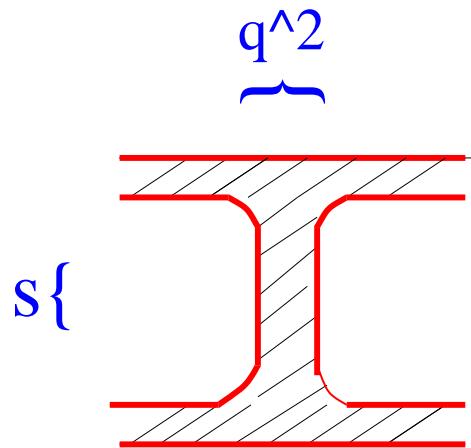
^a(with **Romuald Janik**, Cracow U.) Last papers: hep-th/0512162,
hep-th/0606149

A 35-years-old Puzzle

A String Theory for **Strong** Interactions ?

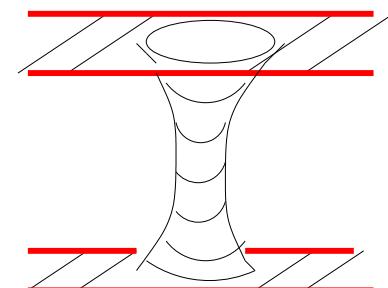
Strings \Rightarrow Gauge Fields \Rightarrow Strings

1968 \Rightarrow 1974 \Rightarrow 1998



Veneziano Amplitude

$$A_R(s, q^2)$$

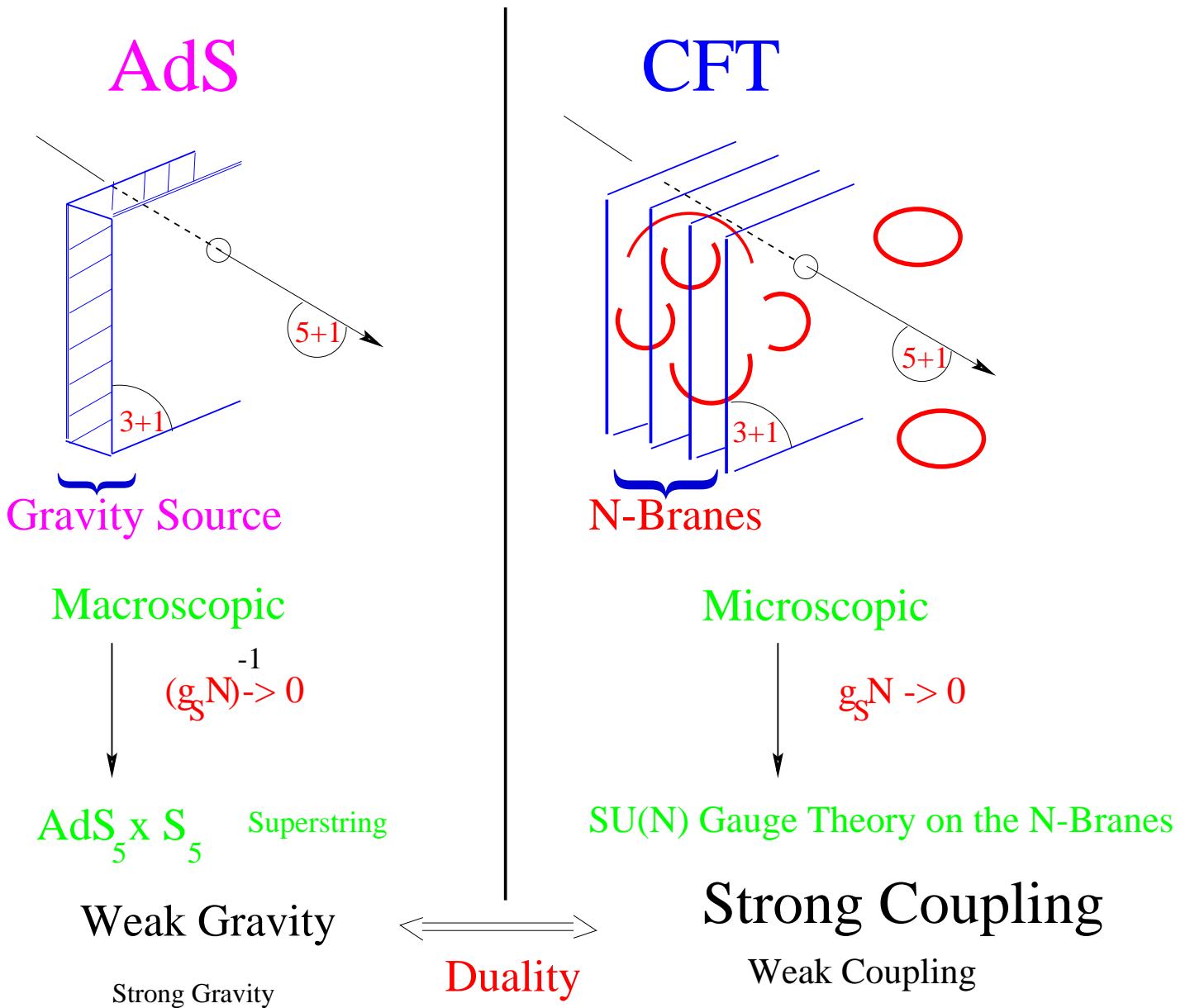


Shapiro-Virasoro Amplitude

$$A_P(s, q^2)$$

AdS/CFT Correspondence

J.Maldacena (1998)



More on AdS/CFT

- **D_3 -brane Solution of Super Gravity:**

$$ds^2 = f^{-1/2}(-dt^2 + \sum_1^3 dx_i^2) + f^{1/2}(dr^2 + r^2 d\Omega_5)$$

“On-Branes \times Out-Branes”

$$f = 1 + \frac{R^4}{r^4} ; R = 4\pi g_{YM}^2 \alpha'^2 N$$

- **“Maldacena limit”:**

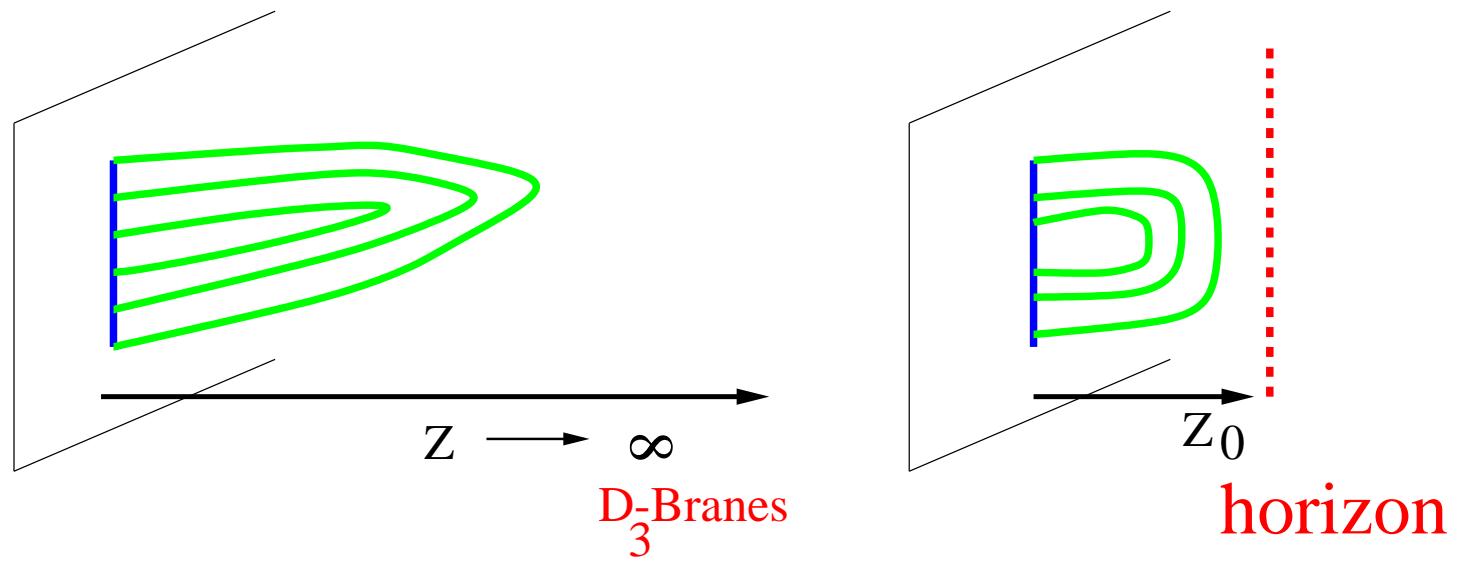
$$\frac{\alpha'(\rightarrow 0)}{r(\rightarrow 0)} \rightarrow z , R \text{ fixed} \Rightarrow g_{YM}^2 N \rightarrow \infty$$

Strong coupling limit

$$ds^2 = \frac{1}{z^2}(-dt^2 + \sum_{1-3} dx_i^2 + dz^2) + R^2 d\Omega_5$$

Background Structure: $\text{AdS}_5 \times S_5$

Holography

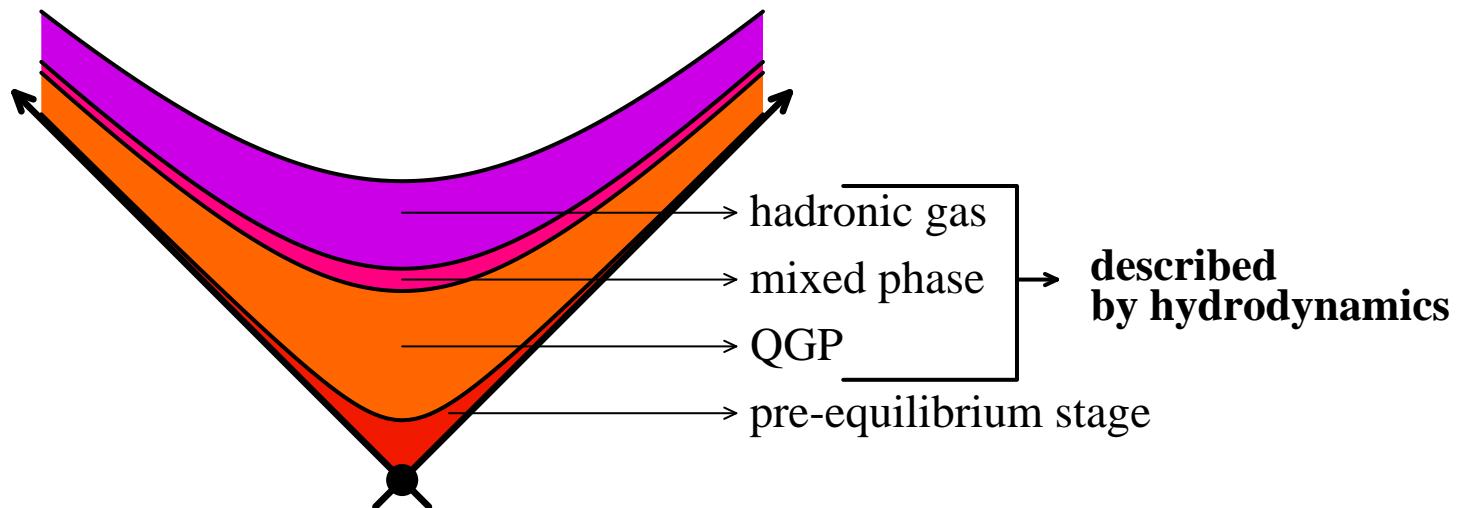


$$Area_{min}^{AdS} = \lim_{\chi \rightarrow \infty} L/(L/\chi) \quad Area_{min}^{BH} = \lim L \times (L/\chi)$$

Confinement or Temperature : Minimal Surfaces with Horizon

AdS/CFT and QCD Hydrodynamics

QGP and Hydrodynamics: J.D.Bjorken (1982)



- Boost Invariance

$$\tau = \sqrt{x_0^2 + x_1^2} ; \quad y = \frac{1}{2} \log \frac{x_0 + x_1}{x_0 - x_1} ; \quad x_T \rightarrow x_1, x_2$$

- QGP: Perfect fluid behaviour
- Pre-equilibrium stage: Fast

The 4d Energy-Momentum Tensor

- Constraint Equations

$$\begin{aligned}\mathbf{T}^{\mu}_{\mu} &\equiv -T_{\tau\tau} + \frac{1}{\tau^2}T_{yy} + 2T_{xx} = 0 \\ \mathbf{D}_{\nu}T^{\mu\nu} &\equiv \tau\frac{d}{d\tau}T_{\tau\tau} + T_{\tau\tau} + \frac{1}{\tau^2}T_{yy} = 0\end{aligned}$$

- Boost-invariant Tensor

$$T_{\mu\nu} = \begin{pmatrix} f(\tau) & 0 & 0 & 0 \\ 0 & -\tau^3 \frac{d}{d\tau} f(\tau) - \tau^2 f(\tau) & 0 & 0 \\ 0 & 0 & f(\tau) + \frac{1}{2}\tau \frac{d}{d\tau} f(\tau) & 0 \\ 0 & 0 & 0 & \dots \end{pmatrix}$$

- 1-parameter family

$$f(\tau) \propto \tau^{-s} : (0 < s < 4 ; T_{\mu\nu}t^{\mu}t^{\nu} \geq 0)$$

$f(\tau) \propto \tau^{-\frac{4}{3}}$: **Perfect Fluid**

$f(\tau) \propto \tau^{-1}$: **Free streaming**

$4 \rightarrow 5d$: Holographic Renormalization

K.Skenderis (2002)

- Fefferman-Graham Coordinates:

$$ds^2 = \frac{g_{\mu\nu}dx^\mu dx^\nu + dz^2}{z^2}$$

- $4d \Leftrightarrow 5d$ metric:

$$g_{\mu\nu} = g_{\mu\nu}^{(0)} (= \eta_{\mu\nu}) + z^2 g_{\mu\nu}^{(2)} (= 0) + z^4 g_{\mu\nu}^{(4)} (\propto \langle T_{\mu\nu} \rangle) + \dots$$

+ ...: from Einstein Eqs.

Technicalities...

- **Boost-Invariant ansatz:**

$$ds^2 = \frac{-e^{a(\tau,z)}d\tau^2 + \tau^2 e^{b(\tau,z)}dy^2 + e^{c(\tau,z)}dx_\perp^2}{z^2} + \frac{dz^2}{z^2}$$

- **Einstein Equation(s):**

$$[a(\tau, z), b(\tau, z), c(\tau, z)] = [a(v), b(v), c(v)] + \mathcal{O}\left(\frac{1}{\tau^\#}\right)$$

$$v = \frac{z}{\tau^{s/4}}$$

$$\begin{aligned} v(2a'(v)c'(v) + a'(v)b'(v) + 2b'(v)c'(v)) - 6a'(v) - 6b'(v) - 12c'(v) + vc'(v)^2 &= 0 \\ 3vc'(v)^2 + vb'(v)^2 + 2vb''(v) + 4vc''(v) - 6b'(v) - 12c'(v) + 2vb'(v)c'(v) &= 0 \\ 2vsb''(v) + 2sb'(v) + 8a'(v) - vsa'(v)b'(v) - 8b'(v) + vsb'(v)^2 + \\ 4vsc''(v) + 4sc'(v) - 2vs a'(v)c'(v) + 2vsc'(v)^2 &= 0 . \end{aligned}$$

- **Asymptotic Solution**

$$\begin{aligned} a(v) &= A(v) - 2m(v) \\ b(v) &= A(v) + (2s - 2)m(v) \\ c(v) &= A(v) + (2 - s)m(v) \end{aligned}$$

(1)

Dual of a Perfect Relativistic fluid

$$v = \frac{z}{\tau^{1/3}}$$

- Asymptotic metric

$$ds^2 = \frac{1}{z^2} \left[-\frac{\left(1 - \frac{e_0}{3} \frac{z^4}{\tau^{4/3}}\right)^2}{1 + \frac{e_0}{3} \frac{z^4}{\tau^{4/3}}} d\tau^2 + \left(1 + \frac{e_0}{3} \frac{z^4}{\tau^{4/3}}\right) (\tau^2 dy^2 + dx_\perp^2) \right] + \frac{dz^2}{z^2}$$

- Black Hole off in the 5th dimension

$$\text{Horizon : } z_0 = \left(\frac{3}{e_0}\right)^{\frac{1}{4}} \cdot \tau^{\frac{1}{3}} .$$

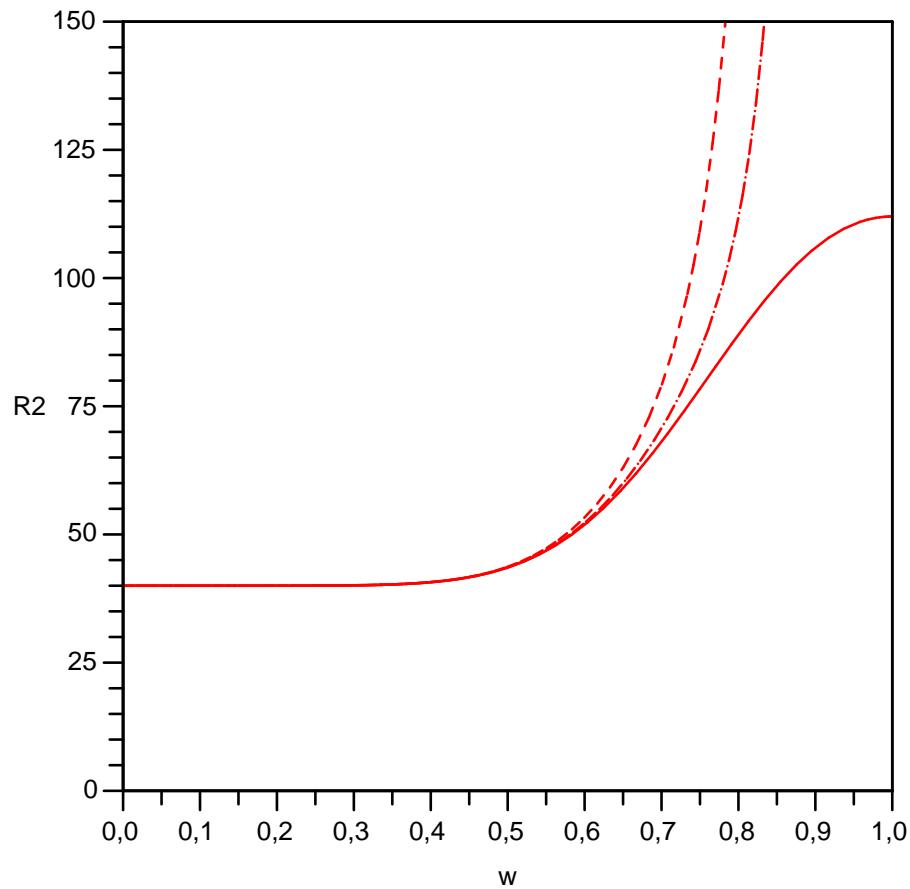
$$\text{Temperature : } T(\tau) \sim \frac{1}{z_0} \sim \tau^{-\frac{1}{3}}$$

$$\text{Entropy : } S(\tau) \sim \text{Area} \sim \tau \cdot \frac{1}{z_0^3} \sim \text{const}$$

- Other cases: True Singularities \neq BH

AdS/CFT: Selection of the Perfect Fluid

$$\mathfrak{R}^2 = R^{\mu\nu\alpha\beta} R_{\mu\nu\alpha\beta} \quad s = \frac{4}{3} \pm .1$$



Thermalization of a perfect fluid

- Thermal stability from scalar excitations:

$$\boxed{\Delta\phi \equiv \frac{1}{\sqrt{-g}} \partial_i (\sqrt{-g} g^{ij} \partial_j \phi) = 0}$$

$$\left[\partial_z \rightarrow \tau^{-\frac{1}{3}} \partial_v ; \quad \partial_\tau \rightarrow \partial_\tau - \frac{1}{3} \tau^{-\frac{4}{3}} \partial_v \right]$$

$$-\frac{1}{v^3} \frac{(1+v^4)^2}{1-v^4} \partial_\tau^2 \phi(\tau, v) + \tau^{-\frac{2}{3}} \partial_v \left(\frac{1}{v^3} (1-v^8) \partial_v \phi(\tau, v) \right) = 0$$

- Scaling solution

$$\partial_\tau^2 f(\tau) = -\omega^2 \tau^{-\frac{2}{3}} f(\tau) \Rightarrow f(\tau) = \sqrt{\tau} J_{\pm\frac{3}{4}} \left(\frac{3}{2} \omega \tau^{\frac{2}{3}} \right) \sim \tau^{\frac{1}{6}} e^{\frac{3}{2} i \omega \tau^{\frac{2}{3}}}$$

$$\boxed{\partial_v \left(\frac{1}{v^3} (1-v^8) \partial_v \phi(v) \right) + \omega^2 \frac{1}{v^3} \frac{(1+v^4)^2}{1-v^4} \phi(v) = 0}$$

- Dominant Decay Frequency

$$\frac{\omega_c}{\pi T} \sim 3.1194 - 2.74667 i$$

Conclusions

- AdS/CFT and Strings
Construction of “Gravity Duals”
- AdS/CFT and QGP Hydrodynamics
Construction of the “Dual” of relativistic fluids
- Selection of the perfect fluid
Non-singular 5d Horizon \Rightarrow Asymptotic
4d perfect fluid
- Thermalization:
The perfect fluid is “very stable”

Unexpected AdS/CFT consequences for QCD at strong coupling?